INFLUENCE OF UNEQUAL COMPONENT EFFICIENCIES ON TRAJECTORIES DURING DISTILLATION OF A QUATERNARY AZEOTROPIC MIXTURE

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ABSTRACT

Experiments were carried out in a bubble cap distillation column operated at total reflux with the system: water (1) – ethanol (2) – methanol (3) – acetone (4). This system has a binary minimum-boiling azeotrope for the water-ethanol mixture and the distillation boundary is represented by a surface with its corners at pure acetone, pure methanol and the water-ethanol azeotrope. For certain starting compositions the measured distillation composition trajectories clearly demonstrate that crossing of the distillation boundary is possible. In order to rationalize our experimental results, we developed a rigorous nonequilibrium (NEQ) stage model, incorporating the Maxwell-Stefan diffusion equations to describe transfer in either fluid phase. The developed NEQ model anticipates the boundary crossing effects, and is in excellent agreement with a series of experiments. In sharp contrast, an equilibrium (EQ) stage model fails even at the *qualitative* level to model the experiments. The differences in the NEQ and EQ trajectories emanates from differences in the component Murphree efficiencies, which in turn can be traced to differences in the binary pair vapour phase diffusivities $\mathcal{D}_{v,ii}$. It is concluded that for reliable design of azeotropic distillation columns we must take interphase mass transfer effects into account in a rigorous manner.

Keywords: quaternary azeotropic distillation, Maxwell-Stefan equations, distillation surface, nonequilibrium stage, equilibrium stage.

INTRODUCTION

Most commercially available simulation programs for distillation columns cater for "real" or non-equilibrium trays [1]. The departure of these real trays from equilibrium behaviour is allowed for in either of two ways. In the first procedure, the user is allowed to specify the individual component Murphree efficiencies for each stage. These component efficiencies can be estimated "off-line" by using the various mass transfer correlations [2,3,4,5,6] as discussed in the standard texts on distillation [1,7,8]. The second approach, which is gaining currency, is to use a fully rate-based approach. In this approach, the interphase mass and heat transfer equations are solved simultaneously along with the interphase equilibrium relations for each stage [9-13]. In the rate-based approaches, the interphase mass transfer relations are invariably based on the Maxwell-Stefan diffusion equations in either fluid phase [13-15]. The rate based approach has been applied in recent times to simulate various complex flow patterns on distillation trays [16,17] and to model maldistribution in packed distillation towers [18]. The rate based approach has also been extended to include three-phase distillation [19], sour-gas absorption [20], and reactive distillation [21].

There is some evidence in the published literature that experimentally measured composition profiles in distillation columns are better simulated with models based on the rigorous Maxwell-Stefan diffusion equations than with simpler models that assume equal component efficiencies [13,14,22-25]. Of particular interest and significance are the experimental measurements, and simulation results, of Pelkonen et al. [24,25]. Pelkonen et al. [24] performed total reflux experiments with the system methanol- iso-propanol -water in a packed distillation column and showed that if the composition at the top of the column is located close to the distillation boundary (i.e. the line connecting pure methanol with the methanol -iso-propanol binary azeotrope) the experimentally measured composition profiles end up with a reboiler composition that is rich in water. The measured composition trajectories can be simulated very well using a nonequilibrium (NEQ) stage model incorporating the Maxwell-Stefan diffusion equations. On the other hand, an equilibrium (EQ) stage model (i.e. a model in which the component efficiencies are each taken to 100%) predicts that the reboiler composition corresponds to pure iso-propanol. The simulation results of Baur et al. [26] show that the differences in the component efficiencies cause the deviation in the NEQ and EQ column trajectories. Pelkonen e tal. [24,25] also performed similar experiments with the quaternary system acetone- methanol- isopropanol- water, with the composition near the top of the column chosen to lie on the distillation boundary, and obtained the same dramatic differences between the predictions of the EQ and NEQ models. The NEQ model predictions were in accord with experiments.

The experimental results of Pelkonen [24] raise the question whether the observed dramatic differences between NEQ and EQ model predictions are also obtained when the starting compositions are not located precisely on the distillation boundary but on either side of it. The major objective of this paper is to demonstrate that for *quaternary* azeotropric distillation, distillation boundaries (which are surfaces dividing the compositions are located within a finite region of compositions on one side of

the boundary. Furthermore, we aim to show that such boundary crossing phenomena can be predicted by the NEQ models, incorporating the Maxwell-Stefan equations and can be attributed to differences in component Murphree efficiencies. Clearly, the EQ models will be incapable of anticipating boundary crossing effects since the EQ distillation trajectories must necessarily follow the residue curve lines for total reflux operations [8,27].

To verify the boundary crossing phenomena, we performed experiments with the quaternary system: water (1) - ethanol (2) - methanol (3) - acetone (4) in a bubble cap tray distillation column. The vapour-liquid equilibrium was determined using NRTL parameters taken from ref [28] and listed in Table 1. The distillation boundary forms a surface connecting the ethanol-water azeotrope with pure methanol and pure acetone. The distillation boundary (surface) is shown in Figs 1 (a), (b) and (c), that represent three different views of the quaternary composition space. Consider Figs 1 (a) and (b), the three dimensional composition space is viewed from two different sides; Fig. 1 (a) shows the component methanol in front whereas Fig. 1 (b) shows the component acetone in front. If only the front of these two projections is considered (the component at the back is set to zero), the projections of the ternary systems: water-ethanol-methanol (Fig 1 (d)) and water-ethanol-acetone (Fig. 1 (e)) remain respectively with their own distillation boundaries (called: the "methanol boundary" in 1 (d) and the "acetone boundary" in 1 (e)). In Fig 1 (f), the two distillation boundaries (Fig. 1 (d-e)) are combined together with all the "pseudo distillation boundary-lines" that lie in between (represented by the gray shaded region); only when a point is located below this shaded region, you can be sure that the point is actually lying below the distillation boundary surface. Fig. 1 (g) shows the same graph as Fig. 1 (f), but with different axis-arrangement.

Component I	Component j	<i>B</i> _{ij} / [K]	<i>B</i> _{ji} / [K]	α _{ij} / [-]
Water	Ethanol	624.92	-29.17	0.294
Water	Methanol	594.63	-182.61	0.297
Water	Acetone	602.62	330.48	0.510
Ethanol	Methanol	73.41	-79.17	0.303
Ethanol	Acetone	188.90	22.83	0.301
Methanol	Acetone	97.78	107.83	0.301

Table 1. NRTL parameters for binary mixtures at 101.3 kPa, taken from Gmehling & Onken [28]. These parameters are used along with $G_{ii} = \exp(-\alpha_{ii}\tau_{ii})$ and $\tau_{ij} = B_{ij}/T$.



Figure 1. (a-b-c). Three dimensional residue curve space for the Water (1) – Ethanol (2) – Methanol (3) – Acetone (4) system, showing an almost plane distillation boundary-surface with its corners at pure methanol, pure acetone and the binary azeotrope between Water-Ethanol. (d) Front view projection of Fig. 2 (a) with the methanol component in front, showing the ternary "methanol boundary". (e) Front view projection of Fig. 2 (b) with the acetone component in front, showing the ternary "acetone boundary. (f-g) Combination of the two ternary boundaries from Fig. 2 (d-e) together with all the "pseudo distillation boundary-lines" that lie in between (represented by the gray shaded region).

EXPERIMENTAL SET-UP

The experiments were carried out in a laboratory-scale distillation column supplied by Schott Nederland B.V.; see Fig. 2. The double layered glass column with vacuum between the inner and outer shell contains a total condenser (stage 1), a partial reboiler (stage 12) and ten equal bubble cap trays (stages 2 to 11) for which the dimensions are tabulated in Table 2 and pictured in the inset to Fig. 2. The distillation column is divided into two sets of five bubble cap trays by an intersection at which a continuous feed can be introduced to the column. Product streams can be tapped automatically from the condenser and manually from the reboiler. The glass distillation column has several small openings of 10 mm in diameter, which are sealed with Teflon-coated septums. These opening enable

liquid and vapour samples to be withdrawn by means of a syringe. The column has a total height of 2160 mm and a 50 mm inner diameter.

The reboiler is placed in a heating mantle, which is connected with a PC provided with the required software (Honeywell: WinNT-workstation 4.0; FIX MMI V 6.15/75-I/O-points runtime; OPTO CONTROL rel.2.2a). By means of the PC, the reboiler temperature can be controlled as well as the feed- and product- flows. Furthermore it provides an automatic safety shut down in case the column reboiler accidentally tends to dry up. The condenser is connected with a water tap, which supplies cooling water to the glass cooling tubes inside the condenser.

Experiments under total reflux conditions and atmospheric pressure were carried out with the quaternary system water – ethanol – methanol - acetone. For any given experiment, 9 vapour and 4 liquid samples were taken from several stages (See Fig. 2) and the temperature profile was measured with Pt 100 sensors. Each sample volume was intentionally kept small (100 μ L) to prevent changes in the composition-profile during the entire experiment. The samples were first dissolved into a reference solvent consisting of 1 vol% cyclohexane in 99 vol% n-propanol before injection into the Gas Chromatograph (type: GC8000-Top with pressure/flow control) by means of an autosampler (type AS800). The channel inside the GC is made of stainless steel and has a total length of 1 m and 0.125 inch diameter. The carrier gas used was Helium because of its high thermal conductivity and chemical inertness. By analyzing samples of pre-prepared, known, compositions, the GC was carefully calibrated. More detailed descriptions of the experimental set-up, measurement technique, GC analysis and composition determination, including pictures of the column and bubble cap trays are available on our web-site: http://ct-cr4.chem.uva.nl/distillation/.



Figure 2. Schematic of laboratory-scale distillation column. Includes total condenser (1), partial reboiler (12), ten bubble cap trays (2-11) and 13 draw-off faucets, 9 for vapor samples (V) and 4 for liquid samples (L). The details of the bubble cap are shown on the right side of the column.

Column diameter	0.0500 m	Hole pitch	0.0142 m
Tray spacing	0.0462 m	Cap diameter	0.0281 m
Number of flow passes	1	Skirt clearance	0.0030 m
Liquid flow path length	0.0308 m	Slot height	0.0050 m
Downcomer clearance	0.0039 m	Active area (of total area)	97.30 %
Deck thickness	0.0030 m	Total hole area (of total area)	8.27 %
Hole diameter	0.0142 m	Downcomer area (of total area)	1.35 %
Weir type	Circular	Slot area	0.000221 m ²
Weir length	0.0182 m	Riser area	0.000158 m ²
Weir height	0.0092 m	Annular area	0.000462 m ²
Weir diameter	0.0058 m		

Table 2. Bubble cap tray design of the laboratory-scale distillation column



Figure 3. (a) – (c) are the results of experiments Q1-01 – Q1-03 (open circles for vapor samples and open squares for liquid samples) showing the column composition trajectories for the Water (1) – Ethanol (2) – Methanol (3) – Acetone (4) system in three different front view projections. Also shown are the simulation results showing the trajectories calculated by the equilibrium (EQ) stage model and the nonequilibrium (NEQ) stage model. The large open circles represent the experimental composition used as input in the simulations. In the NEQ model simulations a bubble size $d_b = 5.0$ mm was chosen.



Figure 3. (d) – (f) are the results of experiments Q1-04 – Q1-06 (open circles for vapor samples and open squares for liquid samples) showing the column composition trajectories for the Water (1) – Ethanol (2) – Methanol (3) – Acetone (4) system in three different front view projections. Also shown are the simulation results showing the trajectories calculated by the equilibrium (EQ) stage model and the nonequilibrium (NEQ) stage model. The large open circles represent the experimental composition used as input in the simulations. In the NEQ model simulations a bubble size $d_b = 5.0$ mm was chosen.

EXPERIMENTAL RESULTS

The experimentally determined composition trajectories for a set of 6 experiments are presented in Fig. 3 (a-f) in three different projections for each single experiment. The first two projections of Fig. 3 (a-f) are similar to the projections shown in Fig. 1 (d) and 1 (e). The third projection is obtained when looked at the distillation boundary surface from above with the component ethanol at the rear (as in Fig 1 (c)). This projection does not give any information concerning a possible boundary crossing, but does show from another point of view the differences between the trajectories predicted by the EQ model versus NEQ model, to be discussed and developed below. At total reflux the composition of the vapor leaving any given stage equals the composition of the liquid arriving at that stage from above. Therefore, the 9 vapor and 4 liquid composition samples can be combined when plotting the composition trajectories. In Fig. 3, the vapor samples are denoted by open circles and the liquid samples by open squares. In experiment Q1-01, the column trajectory was located completely below the distillation boundary surface, which corresponds to the left of the ternary distillation boundaries in Fig. 3 (a-1) and 3 (a-2) (indicated by a thick line). All the remaining five experiments, Q1-02 to Q1-06, clearly exhibit boundary crossing phenomena. Clearly, boundary crossing phenomena is not in conformity with the assumption of thermodynamic phase equilibrium; this is evidenced by the fact that the experimental trajectories do not agree with the constraints, enforced by the distillation boundary (surface). In order to understand, and rationalize, the boundary crossing phenomena we develop a rigorous nonequilibrium (NEQ) stage model.

NONEQUILIBRIUM STAGE MODEL DEVELOPMENT

The development of the NEQ stage model follows the ideas and concepts developed earlier by Taylor, Krishna and others and described in earlier publications [9-15]. A brief review of the model development is given below. Consider first a single stage pictured in Fig. 4.



Figure 4. Schematic of the bubble froth regime on the tray.

All our experiments were carried out in the bubbly flow regime. Visual observations of tray operation show that the bubbles are roughly of uniform size and shape. In our model development we assume that the bubbles rise through the liquid in a plug flow manner. Furthermore, we assume that the liquid phase is well-mixed. The steady state component molar balance for 4-component distillation in tray columns is given by the 3-dimensional matrix relation

$$V_{b} \frac{d(y)}{dh} = [K_{Oy}](y^{*} - y)a'$$
(1)

where *a*' is the interfacial area per unit volume of the dispersed bubble phase and $V_{\rm b}$ is the bubble rise velocity. Equation (1) can be re-written in terms of the overall number of transfer units for the vapor phase [$NTU_{\rm Oy}$]:

$$\frac{dy}{d\xi} = \left[NTU_{Oy}\right] \left(y^* - y\right), \qquad (2)$$

where $\xi = h/h_f$ is the dimensional distance along the froth and $[NTU_{Oy}]$ is defined as:

$$\left[NTU_{Oy}\right] = \int_{0}^{n_{f}} \left[\left[K_{Oy} \right] a' / V_{b} \right] dh$$
(3)

Carrying out the integration, assuming that the matrix of overall mass transfer coefficients $[K_{Oy}]$ does not vary along the froth height, we obtain

$$[NTU_{Oy}] \equiv [K_{Oy}]a'h_f / V_b \equiv [K_{Oy}]a'\tau_V, \qquad (4)$$

From eq. (4), we see that $[NTU_{Oy}]$ can be calculated from knowledge of $[K_{Oy}]$, the interfacial area per unit volume of vapor *a*' and the vapor phase residence time τ_V . In our model we assume that all the bubbles to be spherical in shape with a diameter d_b . The interfacial area per unit volume of vapor *a*' is therefore given by:

$$a' = \frac{6}{d_b} \tag{5}$$

The vapor residence time is determined by:

$$\tau_V = \frac{h_f}{V_b},\tag{6}$$

where $h_{\rm f}$ is the height of dispersion (froth); this is taken to be the height of the downcomer tube above the tray floor, i.e. 9.2 mm as seen in Fig. 2. The bubble rise velocity $V_{\rm b}$ is estimated using the Mendelson equation [29], recommended by Krishna *et.al.* [30]:

$$V_b = \sqrt{\frac{2\sigma}{\rho_L d_b} + \frac{g d_b}{2}}, \tag{7}$$

The overall matrix of mass transfer coefficients $[K_{Oy}]$ is given by the addition of resistances formula:

$$[K_{Oy}]^{-1} = [k_{y}]^{-1} + \frac{c_{t}^{V}}{c_{t}^{L}} [K_{eq}] [k_{x}]^{-1}, \qquad (8)$$

in which $[K_{eq}]$ represents the diagonal matrix of K-values and $[k_y]$ and $[k_x]$ are the partial transfer coefficient matrices for the vapor and liquid phases respectively.

Let us consider the matrix of the multicomponent vapor mass transfer coefficient $[k_y]$. The nine elements $k_{y,ij}$ can be estimated from the mass transfer coefficients of the constituent binary pairs, $\kappa_{y,ij}$ from:

$$\begin{bmatrix} k_y \end{bmatrix} = \begin{bmatrix} R_y \end{bmatrix}^{-1} \tag{9}$$

where the elements of the matrix of inverse mass transfer coefficients $[R_v]$ is given by

$$R_{y,ii} = \frac{y_i}{\kappa_{y,in}} + \sum_{\substack{k=1\\k\neq i}}^4 \frac{y_k}{\kappa_{y,in}}; \quad R_{y,ij} = -y_i \left(\frac{1}{\kappa_{y,ij}} - \frac{1}{\kappa_{y,in}}\right); \quad i = 1, 2, 3$$
(10)

For each of the binary pairs in the mixture, the $\kappa_{y,ij}$ can be estimated from the following equation for instationary diffusion within a spherical bubble [13]:

$$Sh_{ij} = \frac{\kappa_{y,ij}d_b}{D_{y,ij}} = \frac{2}{3}\pi^2 \left(\frac{\sum_{m=1}^{\infty} \exp\{-m^2 \pi^2 F o_{ij}\}}{\sum_{m=1}^{\infty} \frac{1}{m^2} \exp\{-m^2 \pi^2 F o_{ij}\}} \right),$$
(11)

with *ij* = 12,13,14,23,24,34. For Fourier numbers $Fo_{ij} \equiv 4D_{y,ij}\tau_V/d_b^2$ larger than 0.06, the Sherwood number reduces to the asymptotic value:

$$Sh_{ij} = \frac{2\pi^2}{3} \approx 6.58; \quad ij = 12,13,14,23,24,34$$
 (12)

For this steady-state limit, the binary vapor mass transfer coefficients are given by:

$$\kappa_{y,ij} = \frac{2\pi^2}{3} \frac{D_{y,ij}}{d_b}; \quad ij = 12,13,14,23,24,34$$
(13)

Eq. (13) leads to the important conclusion that $\kappa_{y,ij}$ would have an unity-power dependence on the vapor diffusivity $\mathcal{D}_{y,ij}$, which is in sharp contrast with the square-root dependence for small values of Fo, i.e. small vapor phase residence times.

The matrix of the multicomponent liquid mass transfer coefficient $[k_x]$ can be obtained analogously to eq. (9-10). The binary liquid mass transfer coefficient $\kappa_{x,ij}$ can be obtained from the penetration model:

$$\kappa_{x.ij} = 2\sqrt{\frac{D_{x,ij}}{\pi t_c}}; \quad ij = 12,13,14,23,24,34$$
 (14)

where the contact time of the liquid with gas bubbles, t_c is given by:

$$t_c = \frac{d_b}{V_b} \tag{15}$$

In the above set of model equations, the only unknown parameter is the bubble diameter $d_{\rm b}$. Once the bubble diameter is set, the system of equations can be solved. Substituting eq. (8) in eq. (4) gives us the [$NTU_{\rm Oy}$]. Assuming that the [$NTU_{\rm Oy}$] on a single stage is constant, eq. (2) can be integrated using the boundary conditions

$$\xi = 0 \text{ (inlet to tray)}(y) = (y_E)$$

$$\xi = 1 \text{ (outlet of tray)}(y) = (y_L)$$
(16)

to obtain the compositions leaving the distillation stage (detailed derivations are available in reference [23]):

$$(y^* - y_L) = \exp[-[NTU_{Oy}]](y^* - y_E)$$
 (17)

Introducing the matrix $[Q] = \exp[-[NTU_{OV}]]$, we may re-write eq. (17) in the form

$$(y_L - y_E) = [[I] - [Q]](y^* - y_E),$$
 (18)

where [*I*] is the identity matrix. The limiting case of the EQ stage model is obtained when the mass transfer coefficients in either fluid phase attain large values; [*Q*] reduces in this case to the null matrix and the compositions leaving the tray (y_L) are equal to (y^*), in equilibrium with the liquid leaving the tray.

The material balance relations outlined above need to be solved along with the enthalpy balance relations, as described in Chapter 14 of Taylor and Krishna [13]. The required heat transfer coefficients in the vapour phase are calculated from the heat transfer analog of eq. (11) for the vapour phase Nusselt number. Similarly, the liquid phase heat transfer coefficient is obtained by the application of the penetration model to the liquid phase, analogous to eq. (14). The entire set of material and energy balance equations, along with the interphase mass and energy transfer rate relations are then incorporated into a rigorous stage-to-stage model as described in Chapter 14 of Taylor and Krishna [13]. This chapter contains more exhaustive details of this model including sample calculations for binary and ternary mixtures.

SIMULATION STRATEGY

Simulations of the total reflux experimental runs were carried out using both the equilibrium (EQ) stage model and the rigorous nonequilibrium (NEQ) stage model developed above. The operating pressure for all experiments was 101.3 kPa and the ideal gas law was used. Activity coefficients were calculated using the NRTL interaction parameters, specified in Table 1, and the vapor pressures were calculated

using the Antoine equations. The vapor phase was assumed to be thermodynamically ideal. The column consists of 12 stages, including the total condenser (stage 1) and partial reboiler (stage 12). The reflux flow rate (0.006 mol/s) and the bottom flow rate (0.0 mol/s) were used for specifying the column-operations.

Since the column is operated at total reflux, the reflux flow rate determined the inner flow rates of vapor and liquid phases on each stage. Simulation of total reflux operations is "complicated" by the fact that there is no feed to the column at steady-state. To overcome this problem we specify one of the experimentally determined compositions of the streams leaving or entering a stage as input parameter. The simulated composition profile of the total reflux run is forced to pass through this specified composition. In all the experiments we specified the vapor composition leaving stage 4 in performing the simulations. This "input" composition is indicated by the large open circle in Fig. 3. The entire set of equations system was solved numerically by using the Newton's method. The NEQ implementation is available in the software program *ChemSep*, developed by Taylor [9-13]. Detailed information on *ChemSep* are available at http://www.chemsep.org and in the recent book by Kooijman and Taylor [11]; this book contains details of all thermodynamics and mass transfer models for tray columns that have been implemented into the software.

DETERMINATION OF BUBBLE SIZE

Before the quaternary experiments can be simulated, we need to know the bubble size d_b , which is the only "unknown" in the NEQ model developed above. To determine the bubble size we performed further experiments with binary and ternary mixtures: ethanol – water, water-methanol, water-ethanol-acetone, and water-ethanol-methanol. The stage-to-stage composition profiles for two binary mixtures, for a range of compositions, are well represented by the NEQ model taking the bubble size d_b , = 5 mm; see Figs 5 (a), (b) and (c).



Figure 5. Experimental results showing the column composition trajectories for the Water (1)
– Ethanol (2), Water (1) – Methanol (2), Water (1) – Ethanol (2) – Acetone (3) and Water (1)
– Ethanol (2) - Methanol (3) systems. Also shown are the simulation results showing the trajectories calculated by the (NEQ) stage model with a bubble size d_b = 5.0 mm.

The NEQ simulations for the two binary mixtures were carried out by specifying the vapour composition leaving the reboiler (stage 12) as "input" composition in the simulations.

With the bubble size $d_b = 5$ mm, chosen to "fit" the binary experiments, the experiments with the two ternary mixtures are also well-represented by the NEQ model; see Figs 5 (d) and (e). The NEQ simulations for the ternary mixtures were carried out by specifying the vapour composition leaving stage 4 as "input" composition in the simulations. It is interesting to note that the for the water-ethanol-acetone system, the distillation boundary ("acetone" boundary) is crossed; see Fig. 5 (d). Similarly for the water-ethanol-methanol system, the distillation boundary ("methanol" boundary) is crossed; see Fig. 5 (e).

COMPARISON OF EQ AND NEQ SIMULATIONS WITH QUATERNARY EXPERIMENTS

All 6 guaternary experiments, Q1-01 to Q1-06, were simulated with both the EQ stage model and the rigorous NEQ stage model. Let us consider one of the experiments (Q1-06) in some detail. Fig. 6 (a) compares the EQ model with the experimental results. The large open circle represents the vapour composition leaving stage 4; this is specified in the simulations as "input" composition. We note that while the experimental points cross the distillation boundary (grey coloured surface), the EQ column trajectory does not and remains below the boundary surface. A further point to note is that while the experimental results show that proceeding down the column (in the direction of the reboiler) the compositions get richer in water, the EQ simulations predict that these trays get progressively richer in ethanol. The NEQ model simulations require specification of the bubble diameter. The NEQ simulations were carried out for a range of bubble diameters in the 3 - 5.5mm range. Decreasing the bubble diameter has the effect of increasing the mass transfer coefficient (see eqs (13) - (15)) and makes the NEQ model tend towards the EQ model.



Figure 6. (a) EQ model and (b) NEQ model simulation results compared with the experimental data (open circles for vapor samples and open squares for liquid samples) for run Q1-06. The large open circle is the specified composition for the simulations; this corresponds to the vapor composition leaving stage 4.

To match the EQ trajectory, the bubble size has to be 1.5 mm, or smaller. Conversely, increasing the bubble diameter, decreases the mass transfer coefficient and the NEQ trajectories move away from the EQ trajectory. The best agreement with the experiments is obtained with $d_b = 5.0$ mm. This choice of 5 mm is coincident with the binary and ternary mixtures, considered above. The simulation result for the NEQ model, with $d_b = 5.0$ mm, is plotted for the experimental run Q1-06 in Fig 6 (b). The NEQ trajectory is in very good agreement with the experiment results and is able to reproduce the boundary crossing observed.

The simulation results for the EQ and NEQ model (with $d_b = 5.0$ mm) for all the experimental runs are shown in Fig. 3 (a) – (f) in three different projections, along with the experimental results.

Consider the run Q1-01. For this run no boundary crossing is observed experimentally; see Fig. 3 (a-1), (a-2) and (a-3). Both EQ and NEQ models do not anticipate boundary crossing, although the predictions of the NEQ model are superior to that of the EQ model and in much better agreement with the experimentally measured composition trajectories.

Consider the runs Q1-02 up to and including Q1-06 in Fig. 3. For all these runs we experience boundary crossing and the NEQ model successfully anticipates this phenomenon. In all the cases the EQ model fails to cross the boundary and the EQ trajectory remains on one side of the boundary surface. For all these runs the experimental results show that proceeding down the column (in the direction of the reboiler) the compositions get richer in water. The EQ simulations predict that these trays get progressively richer in ethanol; this is qualitatively different to the experimental observations.

COMPONENT MURPHREE EFFICIENCIES IN QUATERNARY DISTILLATION

We may conclude from the foregoing that boundary crossing is caused by *multicomponent* mass transfer effects. To explain this in some detail we consider run Q1-06. The values of the binary pair vapor diffusivities, $\mathcal{D}_{y,ij}$ for water (1) – ethanol (2) – methanol (3) – acetone (4) are specified in Table 3 for stage 6, along with the corresponding liquid phase coefficients and the matrix of vapour phase transfer units [*NTU*_y] and liquid phase transfer units [*NTU*_x]. The estimated values of the Fourier numbers calculated using

 $Fo_{ij} = \frac{4D_{y,ij}\tau_V}{d_b^2} \tag{19}$

are also given in Table 3 for stage 6, along with the values of the surface tension (σ) and liquid density (ρ_L) that are needed in order to estimate the single bubble rise velocity (V_b) and thus the vapor residence time (τ_V). From Table 3, we see that the Fo values exceed 0.06 in all cases, justifying the use of eq. (12) for estimation of the vapor phase mass transfer coefficients $\kappa_{y,ij}$ of the binary pairs in the mixture; the $\kappa_{y,ij}$ have an unity-power dependence on the vapor diffusivities $\mathcal{D}_{y,ij}$. The vapour phase diffusivities of the three binary pairs are estimated using the Fuller-Schettler-Giddings equation; details of the estimation procedure are to be found in Kooijman and Taylor [11]; this book also specifies the estimation methods for liquid phase diffusivities,

densities and surface tension. By evaluating the individual contributions of the liquid and vapor phases in eq. (8) it can be verified that the mass transfer resistance is predominantly in the vapor phase. The liquid phase resistance contributes less than 10% of the total resistance; this conclusion was found to be valid for all the 6 experimental runs carried out in this study.

Table 3. Physical and transport properties for stage 6 of experiment Q1-06 for the Water (1) – Ethanol (2) – Methanol (3) – Acetone (4) system obtained by NEQ model simulations (bubble diameter = 5.0 mm).

Parameter	units	i-j pair					
		1-2	1-3	1-4	2-3	2-4	3-4
		pair	pair	pair	pair	pair	pair
$D_{y,ij}$	[10 ⁻⁵ m²/s]	2.1	2.72	1.82	1.36	0.908	1.18
$D_{x,ij}$	[10 ⁻⁹ m²/s]	6.07	5.52	4.51	4.08	3.07	3.53
NTU _{y,ij}	[-]	1.49	1.93	1.29	0.966	0.644	0.838
NTU _{x,ij}	[-]	16	15.3	13.8	13.1	11.4	12.2
σ	[N/m]	0.03357					
$ ho_{L}$	[kg/m ³]	771.0					
Vb	[m/s]	0.2049					
$\tau_{\rm V}$	[S]	0.0449					
Fo _{ij}	[-]	0.1509 0.1954 0.1308 0.0977 0.0652 0.0848					

To understand the phenomena of boundary crossing, we consider the component Murphree stage efficiencies, defined by

$$E_{i} = \frac{y_{i,L} - y_{i,E}}{y_{i}^{*} - y_{i,E}}, \quad i = 1, 2, 3, 4$$
(20)

For the EQ model the component efficiencies are all equal to unity. For the NEQ model the component efficiencies will, in general, differ from one another. To illustrate this, we present the calculations of E_i for run Q1-06 in Fig. 7 (a) obtained from NEQ simulations with a bubble diameter of 5.0 mm. It is clear that the component Murphree efficiencies are all different from one another and vary from stage to stage. In particular we note that the methanol efficiency is negative on stage 3. The reason for the negative methanol efficiency is that its constituent driving force is vanishingly small on stage 3 (see Fig. 7 (b)) and therefore its transfer is dictated by the movement of the other three components in the mixture. The origin of the differences in E_i can be traced to the differences in the binary pair vapor diffusivities $\mathcal{D}_{y,ij}$. If the binary $\mathcal{D}_{y,ij}$ were close to one another, the differences in the component efficiencies cause the actual composition trajectory followed on any given stage $(y_{i,L} - y_{i,E})$ to deviate from the trajectory dictated by the equilibrium vector $(y_i^* - y_{i,E})$.



Figure 7. (a) Component efficiencies along the column for the experiment Q1-06 calculated by the NEQ stage model. In the NEQ model simulations a bubble size $d_b = 5.0$ mm was chosen. (b) Methanol driving force along the column.

For various vapor compositions entering any given stage, we have plotted in two ways in Fig. 8 (a) and (b) the actual composition vector $(y_{i,L} - y_{i,E})$, calculated from the NEQ model (taking bubble diameter of 5.0 mm) along with the equilibrium vector $(y_i^* - y_{i,E})$. The angle between the NEQ trajectory (continuous line) and the EQ trajectory (dashed line) increases when the differences in the component efficiencies increase. If all the component efficiencies were equal to one another, the NEQ and EQ trajectories would coincide. We see from Fig. 8 (b), that the NEQ trajectory has a tendency to cut across to the right of the EQ trajectory, precisely as has been observed in the experiments (cf. Fig. 3). It is this tendency to cut towards the right of the composition space that causes boundary (surface) crossing.



Figure 8. (a-b) Calculated direction vectors using the EQ stage model (100% efficiency for all components, denoted by dashed lines) and the NEQ stage model (denoted by continuous lines). In the NEQ model simulations a bubble size $d_b = 5.0$ mm was chosen.

CONCLUSIONS

The following major conclusions can be drawn from the work presented in this paper.

- 1. The measured composition trajectories during distillation of water ethanol methanol acetone under total reflux conditions in a bubble cap distillation column clearly demonstrate that crossing of a distillation boundary (surface) is possible.
- 2. An NEQ stage model is able to model the experimental results. The experimental results agree very well with the developed model in which a bubble size of 5.0 mm is chosen. The NEQ model correctly anticipates boundary crossing.
- 3. An EQ stage model fails to anticipate boundary crossing in any experiment. The EQ model provides a much poorer representation of the column composition trajectories and do not even agree qualitatively with the experimental results. While the experimental trajectory shows that the column gets progressively richer in water as we proceed down to the reboiler, the EQ trajectory predicts that the column gets progressively richer in ethanol; see Fig. 3.
- 4. The differences in the NEQ and EQ trajectories emanates from differences in the component Murphree efficiencies, which in turn can be traced to differences in the binary pair vapor phase diffusivities $\mathcal{D}_{y,ij}$.

The overall conclusion to be drawn from this work is that for reliable simulation of distillation of azeotropic systems exhibiting a distillation boundary, we must adopt a rigorous NEQ stage model. In a theoretical simulation study, Castillo and Towler [31] have shown how the differences in the EQ and NEQ distillation column trajectories could be exploited by the engineer in order to obtain process designs that could not be contemplated if mass transfer effects were ignored, and that some designs based solely on EQ models can become infeasible when mass transfer is considered.

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NOTATION

a' B_{ij} c_i c_t d_b \mathcal{D}_{ij} E_i^{MV}	Interfacial area per unit volume of vapour bubbles, m ² m ⁻³ NRTL parameters; see Table 1, K molar concentration of species <i>i</i> , mol m ⁻³ Mixture molar density, mol m ⁻³ Bubble diameter, m Maxwell-Stefan diffusivity for pair i-j, m ² s ⁻¹ Component Murphree point efficiency, dimensionless
Fo	Fourier number, $Fo \equiv 4D_v \tau_V / d_b^2$, dimensionless
G _{ij} g h h _f k _{ij} [K] [K _{eq}] [K _{Oy}] [NTU _{Oy}] [NTU _{Ox}] n [R] S	NRTL parameters; see Table 1, dimensionless Acceleration due to gravity, m s ⁻² distance along froth height, m Height of dispersion, m Element for matrix of multicomponent mass transfer coefficient, m s ⁻¹ Matrix of multicomponent mass transfer coefficients, m s ⁻¹ diagonal matrix of K-values, dimensionless Matrix of multicomponent overall mass transfer coefficients, m s ⁻¹ Matrix of overall number of vapour phase transfer units, dimensionless Matrix of overall number of liquid phase transfer units, dimensionless number of diffusing species, dimensionless Matrix of inverse mass transfer coefficients, m ⁻¹ s parameter defined in eq. (10), m/s
t _c	liquid-bubble contact time, s
Ť	Temperature, K
V _b	Single bubble rise velocity, m s ⁻¹
X _i V:	Vapour composition for component i, dimensionless
Z _i	mole fraction of component i of the appropriate phase
Greek	
$lpha_{ij}$	Non-randomness parameter in NRTL equation, see Table 1, dimensionless
Kij	Binary Maxwell-Stefan mass transfer coefficients, m s ⁻¹
$ ho_{L}$	Density of the liquid, kg m ⁻³
μ_{L}	Liquid viscosity, Pa s
μ_{i}	Surface tension N m ⁻¹
TV	Vapour phase residence time, s
$ au_{ij}$	NRTL parameters; see Table 1, dimensionless
ξ.	Dimensionless distance along dispersion or column height, dimensionless
Subscrint	
b	Referring to a bubble
С	Peterring to conditions entering a choosified stage

- Referring to conditions entering a specified stage Referring to the froth E f

i	Component number
j	Component number
n	Component number
L	Referring to conditions leaving a specified stage
Oy	Overall parameter referred to the vapour phase
ref	Reference
t	Referring to total mixture
х	Referring to the x phase (liquid)
у	Referring to the y phase (vapour)

Superscript

Μ		Referring	to	Murphree

- L Referring to the liquid phase
- V Referring to the vapour phase
- * Referring to equilibrium state

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